

UNIVERSIDAD COMPLUTENSE MADRID



Inflation, Electroweak Vacuum Stability and Direct Searches at LHC

Mindaugas Karčiauskas

Enqvist, MK, Lebedev, Rusak, Zatta (2016) Ema, MK, Lebedev, Zatta (2017) Ema, MK, Lebedev, Rusak, Zatta (2019)

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24 \, {
m GeV}$ Aad et al. (2015)

(Meta) stability region: $111 \text{ GeV} < M_h < 175 \text{ GeV}$

2. No other states discovered

Mindaugas Karčiauskas

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24 \text{ GeV}$ Aad et al. (2015) (Meta) stability region: $111 \text{ GeV} < M_h < 175 \text{ GeV}$



Forum merastability In a an in the case In a an in the case In a an intervention of the University

Mindaugas Karčiauskas

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24 \text{ GeV}$ Aad et al. (2015) (Meta) stability region: $111 \text{ GeV} < M_h < 175 \text{ GeV}$

2. No other states discovered



Tanabashi et al. (2018)

- EW vacuum metastability
 - $\lambda = 0$ at $\sim 10^{10} \, {
 m GeV}$

•
$$V_{\rm max}^{1/4} \sim 10^9 \,{\rm GeV}$$

Lifetime \gg age of the Universe

Mindaugas Karčiauskas

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24$ GeV Aad et al. (2015) (Meta) stability region: 111 GeV $< M_h < 175$ GeV

2. No other states discovered



- EW vacuum metastability
 - $\lambda = 0$ at $\sim 10^{10} \, {
 m GeV}$

•
$$V_{\rm max}^{1/4} \sim 10^9 \,{\rm GeV}$$

– Lifetime \gg age of the Universe

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24$ GeV Aad et al. (2015) (Meta) stability region: 111 GeV $< M_h < 175$ GeV

2. No other states discovered

Markkanen et al. (2018)



EW vacuum metastability

• $\lambda = 0$ at $\sim 10^{10} \, {
m GeV}$

•
$$V_{\rm max}^{1/4} \sim 10^9 \,{\rm GeV}$$

Lifetime \gg age of the Universe

Mindaugas Karčiauskas

1. SM is consistent up to $m_{\rm Pl}$:

Higgs mass: $M_h = 125.09 \pm 0.24$ GeV Aad et al. (2015) (Meta) stability region: 111 GeV $< M_h < 175$ GeV

2. No other states discovered



- EW vacuum metastability
 - $\lambda = 0$ at $\sim 10^{10}$ GeV
 - $V_{\rm max}^{1/4} \sim 10^9 \,{\rm GeV}$
- Lifetime \gg age of the Universe

Questions:

1. Why isn't Higgs in an energetically favourable state? 2. How EW vacuum survived inflation? (If $r > 10^{-4}$)

Mindaugas Karčiauskas

- Answers the two questions
- Minimal: SM + GR + inflaton
- Stable up to the unitarity breaking scale (Planck scale ?)
- Has possible implication for LHC new physics searches

- Generally expect all operators in a Lagrangian that are
 - Lorentz invariant
 - gauge invariant
 - up to dimension 4

- Generally expect all operators in a Lagrangian that are
 - Lorentz invariant
 - gauge invariant
 - up to dimension 4

- In the SM the Higgs bilinear H[†]H is special
 - Lorentz invariant
 - gauge invariant
 - dimension 2

- Generally expect all operators in a Lagrangian that are
 - Lorentz invariant
 - gauge invariant
 - up to dimension 4

- In the SM the Higgs bilinear H[†]H is special
 - Lorentz invariant
 - gauge invariant
 - dimension 2

 $H^\dagger H \phi^2$ and $H^\dagger H \phi$

Mindaugas Karčiauskas

- Generally expect all operators in a Lagrangian that are
 - Lorentz invariant
 - gauge invariant
 - up to dimension 4

- In the SM the Higgs bilinear H[†]H is special
 - Lorentz invariant
 - gauge invariant
 - dimension 2

 $H^{\dagger}H\phi^2$ and $H^{\dagger}H\phi$

A coupling to gravity

 $H^{\dagger}H\hat{R}, \phi\hat{R}, \phi^{2}\hat{R}$

Mindaugas Karčiauskas

- Generally expect all operators in a Lagrangian that are
 - Lorentz invariant
 - gauge invariant
 - up to dimension 4

- In the SM the Higgs bilinear H[†]H is special
 - Lorentz invariant
 - gauge invariant
 - dimension 2

 $H^{\dagger}H\phi^2$ and $H^{\dagger}H\phi$

A coupling to gravity

 $H^{\dagger}H\hat{R}, \phi\hat{R}, \phi^{2}\hat{R}$

Mindaugas Karčiauskas

Mindaugas Karčiauskas

The Model $S = \int \mathrm{d}^4 x \sqrt{-\hat{g}} \left[rac{1}{2} \Omega^2 \hat{R} - rac{1}{2} \hat{g}^{\mu u} \partial_\mu \phi \partial_ u \phi - rac{1}{2} \hat{g}^{\mu u} \partial_\mu h \partial_ u h - V\left(\phi, h ight) ight]$ $\Omega^2 = 1 + \xi_{\phi} \phi^2 + \xi_h h^2$

Mindaugas Karčiauskas

$$V(\phi, h) = \frac{1}{4}\lambda_{h}h^{4} - \frac{1}{2}\mu_{h}^{2}h^{2}$$
$$+ \frac{1}{4}\lambda_{\phi}\phi^{4} + \frac{1}{3}b_{3}\phi^{3} - \frac{1}{2}\mu_{\phi}^{2}\phi^{2} + b_{1}\phi$$
$$+ \frac{1}{2}\lambda_{h\phi}h^{2}\phi^{2} + \sigma h^{2}\phi$$

Mindaugas Karčiauskas



Mindaugas Karčiauskas



Mindaugas Karčiauskas



Mindaugas Karčiauskas



Mindaugas Karčiauskas



Mindaugas Karčiauskas

• $\xi_{\phi} \gg |\xi_h|, 1$

• Dimensional constants $\ll m_{
m Pl}$

- $\xi_{\phi} \gg |\xi_h|, 1$
- Dimensional constants $\ll m_{
 m Pl}$
- Unitarity above inflation scale
- CMB normalisation: $\lambda_{\phi}/\xi_{\phi}^2 \simeq 5 \times 10^{-10}$
- Unitarity cutoff: $\sim \xi_{\phi}^{-1}$

 $\frac{\lambda_{\phi} \left(\Lambda_{I} \right)}{\xi_{\phi} \left(\Lambda_{I} \right)} \lesssim 2 \times 10^{-5}$ $\xi_{\phi} \left(\Lambda_{I} \right) \lesssim 2 \times 10^{2}$

- $\xi_{\phi} \gg \left|\xi_{h}\right|, 1$
- Dimensional constants $\ll m_{
 m Pl}$
- Unitarity above inflation scale
- CMB normalisation: $\lambda_{\phi}/\xi_{\phi}^2 \simeq 5 \times 10^{-10}$
- Unitarity cutoff: $\sim \xi_{\phi}^{-1}$
- No significant radiative corrections

$$\begin{split} \overline{\lambda_{\phi} \left(\Lambda_{I} \right)} &\lesssim 2 \times 10^{-5} \\ \xi_{\phi} \left(\Lambda_{I} \right) &\lesssim 2 \times 10^{2} \\ \lambda_{h\phi} < 10^{-2} \end{split}$$

- $\xi_{\phi} \gg \left|\xi_{h}\right|, 1$
- Dimensional constants $\ll m_{
 m Pl}$
- Unitarity above inflation scale
- CMB normalisation: $\lambda_{\phi}/\xi_{\phi}^2 \simeq 5 imes 10^{-10}$
- Unitarity cutoff: $\sim \xi_{\phi}^{-1}$
- No significant radiative corrections

$$\begin{split} \lambda_{\phi} \left(\Lambda_{I} \right) &\lesssim 2 \times 10^{-5} \\ \xi_{\phi} \left(\Lambda_{I} \right) &\lesssim 2 \times 10^{2} \\ \lambda_{h\phi} &< 10^{-2} \end{split}$$

- n_s agrees with Planck and $r \simeq 3 \times 10^{-3}$
 - potentially observable



Mindaugas Karčiauskas

Ade et al. (2016)

Higgs During Inflation

• Initially $h \sim 0.1 \ m_{\rm Pl}$



Mindaugas Karčiauskas

$$\phi_{\text{end}} \simeq \frac{1}{\sqrt{\xi_{\phi}}}$$



Mindaugas Karčiauskas

$$\phi_{\text{end}} \simeq \frac{1}{\sqrt{\xi_{\phi}}}$$



Mindaugas Karčiauskas

$$\phi_{\rm end} \simeq \frac{1}{\sqrt{\xi_{\phi}}}$$



Mindaugas Karčiauskas

$$\phi_{\rm end} \simeq \frac{1}{\sqrt{\xi_{\phi}}}$$



Mindaugas Karčiauskas

 $\phi_{\text{end}} \simeq \frac{1}{\sqrt{\xi_{\phi}}}$



Mindaugas Karčiauskas

Preheating



Mindaugas Karčiauskas

Preheating

 $\mathcal{L} \supset \sigma h^2 \phi$



Mindaugas Karčiauskas

Reheating

- Reheating temperature $T_{\rm reh} \sim 10^{12} {
 m GeV}$
- Trilinear interaction $\sigma h^2 \phi$ is essential

$$\phi \rightarrow hh$$





Mindaugas Karčiauskas



Mindaugas Karčiauskas



Mindaugas Karčiauskas

• R can be neglected

 $V(\phi, h) = \frac{1}{4}\lambda_{h}h^{4} - \frac{1}{2}\mu_{h}^{2}h^{2} + \frac{1}{4}\lambda_{\phi}\phi^{4} + \frac{1}{3}b_{3}\phi^{3} - \frac{1}{2}\mu_{\phi}^{2}\phi^{2} + b_{1}\phi + \frac{1}{2}\lambda_{h\phi}h^{2}\phi^{2} + \sigma h^{2}\phi$

Mindaugas Karčiauskas

- R can be neglected
- Rotate to mass eigenstates

$$\left(egin{array}{cc} h_1 \ h_2 \end{array}
ight) &=& \left(egin{array}{cc} \cosartheta & \sinartheta \ -\sinartheta & \cosartheta \end{array}
ight) \left(egin{array}{cc} h-v \ \phi \end{array}
ight)$$

- *R* can be neglected
- Rotate to mass eigenstates

$$2\lambda_h v^2 = m_1^2 \cos^2 \vartheta + m_2^2 \sin^2 \vartheta$$
$$\lambda_{h\phi} v^2 - \mu_{\phi}^2 = m_1^2 \sin^2 \vartheta + m_2^2 \cos^2 \vartheta$$
$$\sigma v = \frac{1}{4} \sin 2\vartheta \left(m_1^2 - m_2^2\right)$$

- For $artheta > 0 \; \lambda_h$ is larger than SM \Rightarrow no metastability
- But m_2 cannot be very large \Rightarrow no Landau pole

Mindaugas Karčiauskas

- R can be neglected
- Rotate to mass eigenstates

$$2\lambda_h v^2 = m_1^2 \cos^2 \vartheta + m_2^2 \sin^2 \vartheta$$
$$\lambda_{h\phi} v^2 - \mu_{\phi}^2 = m_1^2 \sin^2 \vartheta + m_2^2 \cos^2 \vartheta$$
$$\sigma v = \frac{1}{4} \sin 2\vartheta \left(m_1^2 - m_2^2 \right)$$



Mindaugas Karčiauskas

- *R* can be neglected
- Rotate to mass eigenstates

$$2\lambda_h v^2 = m_1^2 \cos^2 \vartheta + m_2^2 \sin^2 \vartheta$$
$$\lambda_{h\phi} v^2 - \mu_{\phi}^2 = m_1^2 \sin^2 \vartheta + m_2^2 \cos^2 \vartheta$$
$$\sigma v = \frac{1}{4} \sin 2\vartheta \left(m_1^2 - m_2^2\right)$$

- For $artheta > 0 \; \lambda_h$ is larger than SM \Rightarrow no metastability
- But m_2 cannot be very large \Rightarrow no Landau pole

Mindaugas Karčiauskas

- *R* can be neglected
- Rotate to mass eigenstates

$$2\lambda_h v^2 = m_1^2 \cos^2 \vartheta + m_2^2 \sin^2 \vartheta$$
$$\lambda_{h\phi} v^2 - \mu_{\phi}^2 = m_1^2 \sin^2 \vartheta + m_2^2 \cos^2 \vartheta$$
$$\sigma v = \frac{1}{4} \sin 2\vartheta \left(m_1^2 - m_2^2\right)$$

- For $artheta > 0 \; \lambda_h$ is larger than SM \Rightarrow no metastability
- But m_2 cannot be very large \Rightarrow no Landau pole

Mindaugas Karčiauskas

RG running

- Perturbativity
- Stability
- EW vacuum is global
- For $\sigma \lesssim$ TeV and $m_2 \lesssim$ TeV stable EW vacuum
- For $\langle \phi \rangle$ up to $\sim 10 \,\text{TeV}$
 - Mixing is governed by σ





Mindaugas Karčiauskas

Prospects for LHC

Universal reduction in the Higgs coupling to gauge bosons and fermions

$$\mathcal{L} \supset \frac{H_1 \cos \theta + H_2 \sin \theta}{v} \left[2m_W^2 W_\mu^+ W^{\mu-} + m_Z^2 Z_\mu Z^\mu - \sum m_f \bar{f} f \right]$$

$$\downarrow \\ \sigma \left(pp \to h_1 \right) = \cos^2 \theta \sigma_{\text{SM}} \left(pp \to h_1 \right)$$

Prospects for LHC

Universal reduction in the Higgs coupling to gauge bosons and fermions

$$\mathcal{L} \supset \frac{H_1 \cos \theta + H_2 \sin \theta}{v} \left[2m_W^2 W_\mu^+ W^{\mu-} + m_Z^2 Z_\mu Z^\mu - \sum m_f \bar{f} f \right]$$
$$\underset{\sigma (pp \to h_1) = \cos^2 \theta \sigma_{\text{SM}} (pp \to h_1)$$

- Current constraints: $|\sin \theta| < 0.3$ ATLAS: Aad et al. (2015)
- Deviations of $\sim O(\%)$ will be possible to detect at high luminosity LHC
 - That corresponds to $|\sin \theta| < 0.1$
- Heavy Higgs-like resonance: $h_2
 ightarrow b_2$
 - With $|\sinartheta|\sim 0.1$ and $m_2\lesssim$ TeV = accessible to LHC searches

Prospects for LHC

Universal reduction in the Higgs coupling to gauge bosons and fermions

$$\mathcal{L} \supset \frac{H_1 \cos \theta + H_2 \sin \theta}{v} \left[2m_W^2 W_\mu^+ W^{\mu-} + m_Z^2 Z_\mu Z^\mu - \sum m_f \bar{f} f \right]$$
$$\downarrow \\ \sigma \left(pp \to h_1 \right) = \cos^2 \theta \sigma_{\text{SM}} \left(pp \to h_1 \right)$$

- Current constraints: $|\sin \theta| < 0.3$ ATLAS: Aad et al. (2015)
- Deviations of $\sim \mathcal{O}(\%)$ will be possible to detect at high luminosity LHC
 - That corresponds to $|\sin \theta| < 0.1$
- Heavy Higgs-like resonance: $h_2
 ightarrow h_1 h_1$
 - With $|\sin \vartheta| \sim 0.1$ and $m_2 \lesssim \text{TeV} \Rightarrow$ accessible to LHC searches

Lewis & Sullivan (2017)



Summary

A minimal model with generically expected interactions

- The trilinear interaction is essential
- Assume only one scale: dimensional parameters \sim TeV
- Provides a consistent history of the Universe
- Solves EW vacuum metastability problem
- Suggests new signatures for the high luminosity LHC searches

Enqvist, MK, Lebedev, Rusak, Zatta (2016) Ema, MK, Lebedev, Zatta (2017) Ema, MK, Lebedev, Rusak, Zatta (2019)

Mindaugas Karčiauskas

Mindaugas Karčiauskas

Production of Infrared Modes

Espinosa et al. (2015)



Langevin equation

$$\frac{\mathrm{d}h}{\mathrm{d}N} \simeq -\frac{\lambda h^3}{3H^2} + \frac{H}{2\pi}\xi\left(N\right)$$

• Safe energy scale for inflation $V_{
m inf}^{1/4} < 10^{13} \, {
m GeV}$

• If GW are detected ($r>10^{-4}$) $V_{
m inf}^{1/4}~>~10^{15}~{
m GeV}$

Mindaugas Karčiauskas

Field Shift

• Field shift $\phi o \phi' = \phi - u$ transform dimensional constants as

$$\begin{array}{rcl} b_{3}' &=& b_{3} + 3\lambda_{\phi}u\\ {\mu_{\phi}'}^{2} &=& \mu_{\phi}^{2} - 3\lambda_{\phi}u^{2} - 2b_{3}u\\ {\mu_{h}'}^{2} &=& \mu_{h}^{2} - \frac{1}{2}\lambda_{h\phi}u^{2} - 2\sigma u\\ b_{1}' &=& b_{1} + \lambda_{\phi}u^{3} + b_{3}u^{2} - \mu_{\phi}^{2}d\\ \sigma' &=& \sigma + \lambda_{h\phi}u \end{array}$$

where

$$u \equiv \langle \phi \rangle$$

Dimensionless constants do not change

Mindaugas Karčiauskas