# <span id="page-0-0"></span>Observational constraints on a Unified Dark Matter model with fast transition

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- <span id="page-2-0"></span>Unified Dark Matter (UDM): a single matter component explains both structure formation and cosmic acceleration
- UDM are appealing because evade the coincidence problem and predict  $w_{DE} \approx -1$
- Appearance of  $c_{\rm s}^2\neq$  0 (non-negligible Jeans scale/strong late ISW effect)
- Possible solution: UDM models with fast transition between an early CDM-like phase and a late ΛCDM-like epoch
- Our aim: study the observational viability at structure formation level of a UDM model with fast transition

<span id="page-3-0"></span>The energy density of this UDM model (see M. Bruni al, MNRAS 431 (2013) 2907-2916) is given by

$$
\rho = \begin{cases} \rho_{\rm t} \left( \frac{a_{\rm t}}{a} \right)^3 & a < a_{\rm t} \\ \rho_{\Lambda} + \left( \rho_{\rm t} - \rho_{\Lambda} \right) \left( \frac{a_{\rm t}}{a} \right)^3 & a > a_{\rm t} \end{cases} \tag{1}
$$

where  $a_t$  represents the value of the scale factor at the transition and  $\beta$ refers to the rapidity of the transition. In addition to  $a_t$  and  $\beta$ , there is a third parameter in this model,  $\rho_{\Lambda}$  (or equivalently  $\Omega_{\Lambda}$ ), that plays the role of an effective cosmological constant.

It is useful to explicitly incorporate a Heaviside function  $H(a - a_t)$  so that

$$
\rho = \rho_{t} \left(\frac{a_{t}}{a}\right)^{3} + \rho_{\Lambda} \left[1 - \left(\frac{a_{t}}{a}\right)^{3}\right] H(a - a_{t}), \qquad (2)
$$

We shall use the following continuous approximation to the Heaviside function

$$
H_{t}(a-a_{t})=\frac{1}{2}+\frac{1}{\pi}\arctan(\beta(a-a_{t})),
$$
\n(3)

## <span id="page-5-0"></span>Conditions for viability of UDM models

- Any viable UDM model should satisfy the condition  $k_\mathrm{J}^2 \gg k^2$  for all the scales of cosmological interest
- The explicit form of the Jeans wave number is

$$
k_{\rm J}^2 = \frac{3}{2}\rho a^2 \frac{(1+w)}{c_{\rm s}^2} \left| \frac{1}{2}(c_{\rm s}^2 - w) - \rho \frac{dc_{\rm s}^2}{d\rho} + \frac{3(c_{\rm s}^2 - w)^2 - 2(c_{\rm s}^2 - w)}{6(1+w)} + \frac{1}{3} \right| \ . \tag{4}
$$

We can obtain a large  $k_\mathrm{J}^2$  when we have  $c_\mathrm{s}^2=0$  but also when  $c_\mathrm{s}^2$ changes rapidly. This is the motivation to build UDM models with fast transition.

#### Speed of sound in fast transition models



**•** The amplitude depends on  $a_t$  and  $\beta$ . Faster transitions produce higher and narrower peaks. Earlier transitions for equal rapidities produce peaks with smaller areas.

#### Matter power spectra for different *a<sup>t</sup>* and β



Figure: Dependence of the matter power spectrum on  $a_t$  and  $\beta$ .  $a_t$  increases from bottom to top and β increases from right to left, while  $Ω<sub>Δ*u*dm</sub>$  is kept fixed at 0.7. The crosses show the effective scales of the structure formation data used in our analyses

## <span id="page-8-0"></span>KiDS data points used

- When including KiDS data, in principle we need to consider the non-linear matter power spectrum. The problem is that when there are oscillations in it, we cannot apply a non-linear correction (HALOFIT).
- We found that  $k = 0.28 0.84$  is the widest range of the matter power spectrum probed by the KiDS lensing power spectra. These two data points are then in the mildly non-linear regime, since the threshold is usually considered to be  $k = 0.2 h/Mpc$ .
- We computed the matter power spectrum for ΛCDM model with and without HALOFIT. The linear and non-linear power spectra are identical for large scales (*k* < 0.2 h/Mpc) and deviate for smaller scales. We then computed the likelihood of these ΛCDM power spectra using only these two points in the five KiDS power spectra. We found that the two likelihoods (ΛCDM with and without HALOFIT) deviate by 10%.
- Therefore, we are implicitly introducing an extra theoretical uncertainty of around 10% in the statistical analysis.

<span id="page-9-0"></span>For several combinations of  $\beta$ ,  $\Omega_{\Lambda}$ <sub>ωdm</sub>, and  $a_t$ , we find that  $c_{\rm g}^2 > 1$ . In fact, for very high  $\beta$ , only a very small region of  $\boldsymbol{a}_t$  gives a  $c_\text{s}^2 \leq 1$  (e.g.,  $β = 500000$  and  $Ω<sub>Δ*u*dm</sub>$  = 0.72, only the range  $a<sub>t</sub> < 0.041$  gives  $c_s^2 \leq$  1). This means that if we try to run a chain exploring a wide prior for  $a_t$  and  $\beta$  at the same time such as  $a_t \in [0, 1]$  and  $\beta \in [0, 500000]$ , a very large amount of points will be rejected, which could be a problem if there is not enough time and computer resources.



Table: Ranges for the UDM parameters used in the MCMC runs.



Figure: Regime 1: Posterior probabilities for each parameter in the vanilla set-up, as well as the contours with the 1- $\sigma$  and 2- $\sigma$  confidence regions for the analysis with (black lines) and without KiDS (blue lines).



Figure: Regime 2: Posterior probabilities for each parameter in the vanilla set-up, as well as the contours with the 1- $\sigma$  and 2- $\sigma$  confidence regions for the analysis with (black lines) and without KiDS (blue lines).



Figure: Regime 3: Posterior probabilities for each parameter in the vanilla set-up, as well as the contours with the 1- $\sigma$  and 2- $\sigma$  confidence regions for the analysis with (black lines) and without KiDS (blue lines).



 $-$  ln  $\mathcal{L}_{\text{min}} = 5379.71$ , minimum  $\chi^2 = 1.076$ *e* + 04

Table: Estimated best-fit, mean, 1- $\sigma$  uncertainty and the 2- $\sigma$  intervals constraints for regime 3 without KiDS data.



 $-$  ln  $\mathcal{L}_{\text{min}} = 54$ 13.66, minimum  $\chi^2 = 1.083e+04$ 

Table: Estimated best-fit, mean, 1- $\sigma$  uncertainty and the 2- $\sigma$  intervals constraints for regime 3 with KiDS data.

## <span id="page-15-0"></span>Nested sampling analysis



Figure: Posterior probabilities for each parameter in the vanilla set-up, as well as the contours with the 1- $\sigma$  and 2- $\sigma$  confidence regions for the analysis with (blue lines) and without KiDS (red lines)



Table: The ranges for the UDM parameters used in the nested sampling analysis



Table: Estimated best-fit, mean, 1- $\sigma$  uncertainty and the 2- $\sigma$  intervals constraints for the nested sampling analysis without KiDS data.

## Nested sampling analysis



 $-$  ln  $\mathcal{L}_{\text{min}} = 5392.94$ , minimum  $\chi^2 = 1.079e+04$ 

Table: Estimated best-fit, mean, 1- $\sigma$  uncertainty and the 2- $\sigma$  intervals constraints for the nested sampling analysis with KiDS data.

<span id="page-19-0"></span>

Table: Values from model comparison criteria, computed for the MCMC and Nested Sampling analyes of the UDM model and for a ΛCDM control analysis.

The nested sampling analysis shows a very good best-fit for the UDM model that is better than the ΛCDM one, even when reduced by the number of degrees-of-freedom. This means that the penalty of having a larger number of extra parameters, does not prevent it from having comparable numbers to <code>ACDM</code> in the reduced  $\chi^2,$  AIC and BIC criteria.

### <span id="page-20-0"></span>**Conclusions**

- UDM models have the advantage that they can describe the dynamics of the Universe with a single dark fluid which triggers both the accelerated expansion at late times and the LSS formation at earlier times
- UDM models have no coincidence problem by definition and predict an effective cosmological constant at late times
- We have constrained observationally, at structure formation level, a UDM model with a fast transition between an Einstein de Sitter model, and a more recent epoch whose dynamics, background and perturbative, are that of a standard ΛCDM model.
- We have shown that for an early enough fast transition our UDM model is compatible with observations at background, linear and midly non-linear level and that it is a viable alternative to ΛCDM.
- Next: Spherical collapse and N-body.