

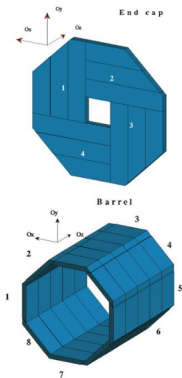
Analysis of the Higgsstrahlung process
($e^+e^- \rightarrow Z(qq)H$) in the hadronic decay mode for the
context of the ILC with the SDHCAL.

Héctor García Cabrera

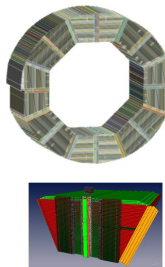
CIEMAT

- Analysis of the Higgsstrahlung process in the context of the ILC.
- The detector model used is the ILD /5_o2_v02. Large model with the SDHCAL in the Tesla geometry and the SiWECal. [List of all models.](#)

SiWECAL



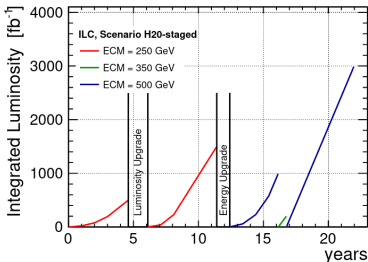
SDHCAL



- Original scenarios in the proposal of 2015: [ILC Operating Scenarios](#)

\sqrt{s}	$\int \mathcal{L} dt$ [fb ⁻¹]			
	G-20	H-20	I-20	Snow
250 GeV	500	2000	500	1150
350 GeV	200	200	1700	200
500 GeV	5000	4000	4000	1600

- H-20 scenario consolidated in the report to Snowmass 2021: [ILC report to Snowmass 2021](#)



- At the moment of the request the request the Snowmass 2015 scenario was used as reference. $L = 1150 \text{ fb} \rightarrow 10.5$ years of operation, which equals to ~ 8.5 years in the H-20 scenario due to the early luminosity upgrade.
- The beam polarization values of operation in all scenarios are P_{e^-} (80%) and P_{e^+} (30%) creating the following proportions in the case of the $(-, +)$ as example:

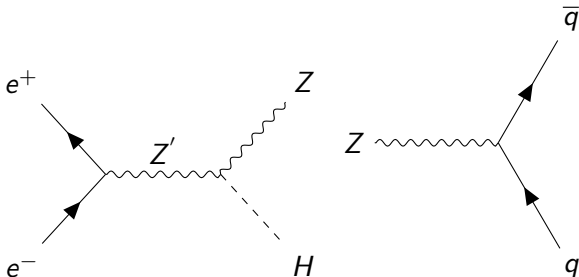
$$\begin{aligned} eL &= 90\% & eR &= 10\% \\ pL &= 35\% & pR &= 65\% \end{aligned}$$

- The Snowmass 2015 scenario has the polarization mixture of: $(-, +) = 67.5\%$, $(+, -) = 22.5\%$, $(-, -) = 5\%$ and $(+, +) = 5\%$. Producing the following polarization factors for the total samples:

$$\begin{aligned} eLpR &= 42.175\% & eLpL &= 25.825\% \\ eRpL &= 17.425\% & eRpR &= 14.575\% \end{aligned}$$

- H-20 polarization mixture: $(-, +) = 45\%$, $(+, -) = 45\%$, $(-, -) = 5\%$ and $(+, +) = 5\%$. Easily adapted from the Snowmass 2021 by weighting the histograms.

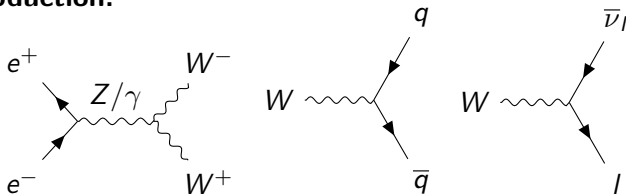
The signal studied in this analysis is the Higgsstrahlung process with the Z decaying hadronically:



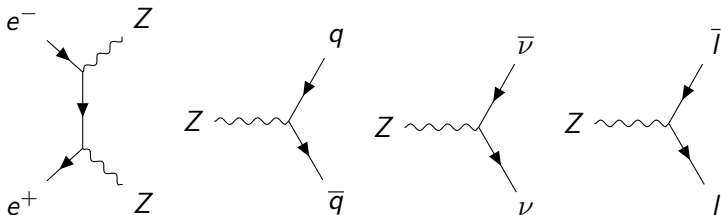
Also the MC request includes all the Higgs decay modes to make per channel tests: $H \rightarrow b\bar{b}$, $H \rightarrow c\bar{c}$, $H \rightarrow gg$, $H \rightarrow WW^*$, $H \rightarrow ZZ^*$, $H \rightarrow \mu\bar{\mu}$, $H \rightarrow \tau\bar{\tau}$, $H \rightarrow \gamma\gamma$ and $H \rightarrow Z\gamma$

The background includes all processes with two or more jets in the final state.

W pair production.

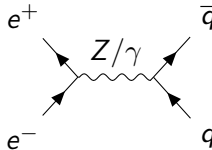


Z pair production.

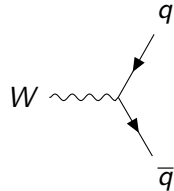
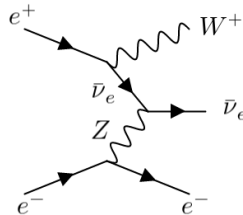
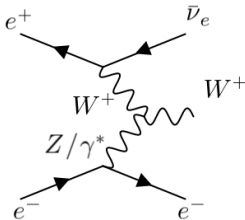


The background includes all processes with two or more jets in the final state.

Z hadronic decay.

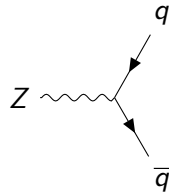
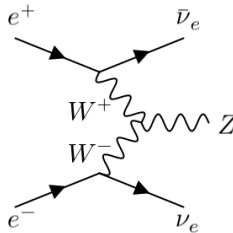
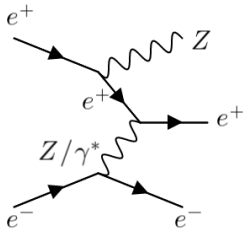


Single W production (+ charge conjugation).



The background includes all processes with two or more jets in the final state.

Single Z production (+ charge conjugation).



Process	eLpR (58.5%)	eRpL (3.5%)	eLpL (31.5%)	eRpR (6.5%)	Size
	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} / NEvts / ReqEvts	
Pqqh	3.4303023e+02 / 211049	2.1948615e+02 / 8834	-	-	28.5 GB
Pqqh_aa	7.78678622e-01 / 479 / 10k	4.98233561e-01 / 20 / -	-	-	1.3 GB
Pqqh_az	5.24836252e-01 / 322 / 10k	3.35813810e-01 / 13 / -	-	-	1.3 GB
Pqqh_bb	1.99643594e+02 / 122830 / 100k	1.27740939e+02 / 5141 / 10k	-	-	14.3 GB
Pqqh_cc	9.91357365e+00 / 6099 / 10k	6.34314973e+00 / 255 / 10k	-	-	2.6 GB
Pqqh_e2e	1.781501e+01 / 1000	1.781501e+01 / 1000	-	-	-
Pqqh_e3e3	2.15079954e+01 / 13232 / 10k	1.37617816e+01 / 553 / 10k	-	-	2.6 GB
Pqqh_gg	2.80941758e+01 / 17284 / 10k	1.79759157e+01 / 723 / 10k	-	-	2.6 GB
Pqqh_ww	7.34084692e+01 / 45164 / 50k	4.69700361e+01 / 1890 / 10k	-	-	7.8 GB
Pqqh_zz	8.98739203e+00 / 5529 / 10k	5.75053713e+00 / 231 / 10k	-	-	2.6 GB
Total					64 GB

Table 1: Full crosssections σ_{Full} , number of events from the Snowmass scenario ($NEvts = L * \sigma_{Full} * f_{Pol}$; $L = 1150$ fb and f_{Pol} being the respective polarization factor), and the requested number of events per event and polarization. The last column is the estimated size of the total requested sample.

Process	eLpR (58.5%)	eRpL (3.5%)	eLpL (31.5%)	eRpR (6.5%)	Size
	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} (fb) / NEvts / ReqEvts	σ_{Full} / NEvts / ReqEvts	
P2f_z_h	1.27965530e+05 / 78730792 / 1M	7.04167430e+04 / 2834273 / 1M	-	-	260 GB
P4f_sw_sl	1.02640160e+04 / 6314935 / 1M	8.66961490e+01 / 3489 / 10k	1.90531440e+02 / 69020 / 100k	1.90637490e+02 / 14250 / 20k	147 GB
P4f_sze_sl	1.42330980e+03 / 875691 / 1M	1.21939670e+03 / 49080 / 50k	1.15583340e+03 / 418700 / 400k	1.15720060e+03 / 86500 / 100k	202 GB
P4f_szm_u_sl	4.53869760e+02 / 279243 / 300k	1.31219580e+02 / 5281 / 10k	-	-	40 GB
P4f_ww_h	1.48664200e+04 / 9146564 / 1M	1.36821530e+02 / 5507 / 10k	-	-	131 GB
P4f_ww_sl	1.87791450e+04 / 11553868 / 1M	1.73468290e+02 / 6982 / 10k	-	-	131 GB
P4f_zz_h	1.40506000e+03 / 864463 / 1M	6.06709780e+02 / 24420 / 30k	-	-	134 GB
P4f_zzm_u_sl	6.09878860e+02 / 375227 / 400k	2.61567200e+02 / 10528 / 10k	-	-	54 GB
P4f_zzorww_h	1.23892920e+04 / 7622511 / 1M	2.25568680e+02 / 9079 / 10k	-	-	131 GB
P4f_zz_sl	8.38079490e+02 / 515628 / 500k	4.66816440e+02 / 18789 / 20k	-	-	68 GB
Total					1.3 TB

Table 2: Full crossections σ_{Full} , number of events from the Snowmass scenario ($NEvts = L * \sigma_{Full} * f_{Pol}$; $L = 1150$ fb and f_{Pol} being the respective polarization factor), and the requested number of events per event and polarization. The last column is the estimated size of the total requested sample.

Process Name	Process ID	Polarization	Requested Events
Pqqh	I402011	cLpR	211049
Pqqh	I402012	cRpL	8834
P2f.z.h	I500010	cLpR	1000000 (1M)
P2f.z.h	I500012	eRpL	1000000 (1M)
P4f_sw_sl	I500105	eLpL	100000 (100k)
P4f_sw_sl	I500106	cLpR	1000000 (1M)
P4f_sw_sl	I500108	cRpL	10000 (10k)
P4f_sw_sl	I500107	eRpR	20000 (20k)
P4f_sze_sl	I500101	eLpL	400000 (400k)
P4f_sze_sl	I500102	cLpR	1000000 (1M)
P4f_sze_sl	I500104	eRpL	50000 (50k)
P4f_sze_sl	I500103	eRpR	100000 (100k)
P4f_szm_sl	I500110	cLpR	300000 (300k)
P4f_szm_sl	I500112	cRpL	10000 (10k)
P4f_ww.h	I500066	cLpR	1000000 (1M)
P4f_ww.h	I500068	eRpL	10000 (10k)
P4f_ww_sl	I500082	cLpR	1000000 (1M)
P4f_ww_sl	I500084	eRpL	10000 (10k)
P4f_zz.h	I500062	cLpR	1000000 (1M)
P4f_zz.h	I500064	eRpL	30000 (30k)
P4f_zzm_sl	I500078	eLpR	400000 (400k)
P4f_zzm_sl	I500080	eRpL	10000 (10k)
P4f_zzorww.h	I500070	cLpR	1000000 (1M)
P4f_zzorww.h	I500072	eRpL	10000 (10k)
P4f_zz_sl	I500074	cLpR	500000 (500k)
P4f_zz_sl	I500076	cRpL	20000 (20k)
Pqqh.aa	I402214	cLpR	10000 (10k)
Pqqh.az	I402215	cLpR	10000 (10k)
Pqqh.bb	I402209	cLpR	100000 (100k)
Pqqh.bb	I402218	eRpL	10000 (10k)
Pqqh.cc	I402210	cLpR	10000 (10k)
Pqqh.cc	I402219	eRpL	10000 (10k)
Pqqh.e3e3	I402216	cLpR	10000 (10k)
Pqqh.e3e3	I402225	eRpL	10000 (10k)
Pqqh.gg	I402211	cLpR	10000 (10k)
Pqqh.gg	I402220	eRpL	10000 (10k)
Pqqh.ww	I402212	cLpR	50000 (50k)
Pqqh.ww	I402221	eRpL	10000 (10k)
Pqqh.zz	I402213	cLpR	10000 (10k)
Pqqh.zz	I402222	eRpL	10000 (10k)

- The samples include the simulation, reconstruction, digitization and overlay.
- The next step is to reconstruct the jets using FastJet.
- An event is selected if it has a dijet compatible with the Z mass.
- The parameters from FastJet are optimized by maximizing the percentage of well reconstructed jets. ($Z_{Tag} > 0.9$ and $Z_{Purity} > 0.9$).

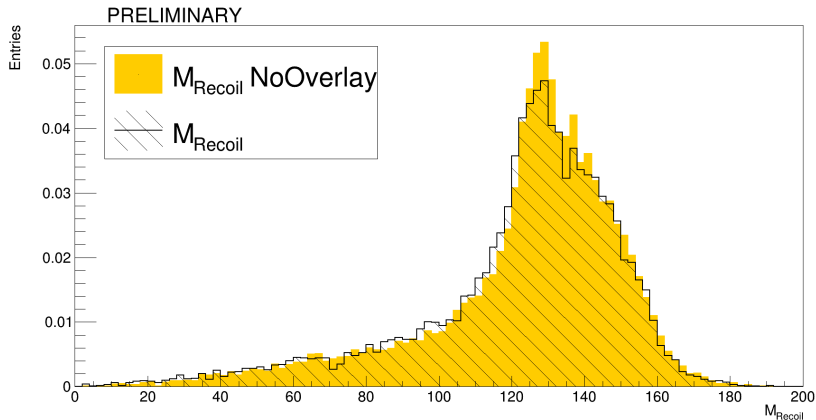
$$\{j_1, j_2\} = \mathit{argmin} |m_{j_1, j_2} - m_Z| ; j_1, j_2 \in N_{Jets} ; j_1 \neq j_2$$

$$Z_{Tag} = \frac{E_Z^{j_1, j_2}}{E_Z^{Total}} ; Z_{Purity} = \frac{E_Z^{j_1, j_2}}{E^{j_1, j_2}}$$

Taking into account the incident angle of the two beams of 14mrad and $\sqrt{s} = 250\text{GeV}$ the computations of the recoil mass are the following:

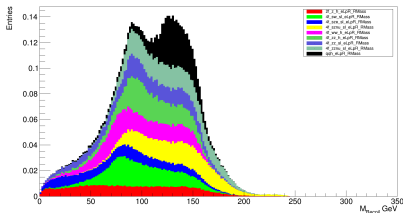
$$E_{q\bar{q}} = E_q + E_{\bar{q}} ; \quad \vec{p}_{q\bar{q}} = \vec{p}_q + \vec{p}_{\bar{q}} ; \quad M_Z = \sqrt{E_{q\bar{q}}^2 - \vec{p}_{q\bar{q}}^2}$$
$$M_{Recoil}^2 = (2E + E_{q\bar{q}})^2 - ((2p_x + p_{q\bar{q}x})^2 + p_{q\bar{q}y}^2 + p_{q\bar{q}z}^2)$$
$$2E = 250.0061252\text{GeV} ; \quad 2p_x = 1.7500286\text{GeV}$$

With this algorithm the optimal parameters have been found to be $Y_{\text{Cut}} = 0.002$ and an impact of less than 0.3 % from overlay in the Z reconstruction in the study of the signal:

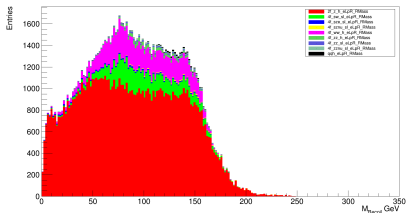


This steps leads to this recoil mass without the application of cuts to the background:

UnWeighted



Weighted



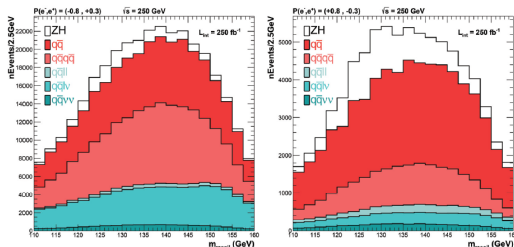
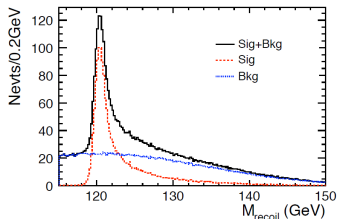
$$\sigma_{ZH} = \frac{N_{Total} - N_B}{BR(Z \rightarrow q\bar{q}) \cdot \epsilon \cdot L}$$

The error of the cross section is computed as:

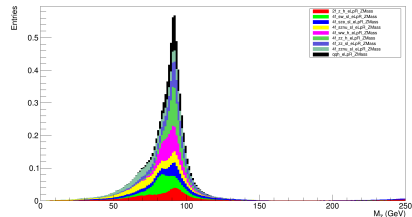
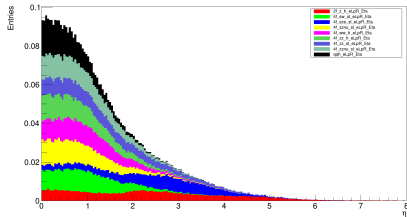
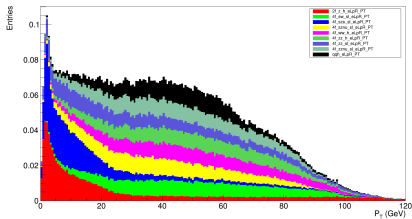
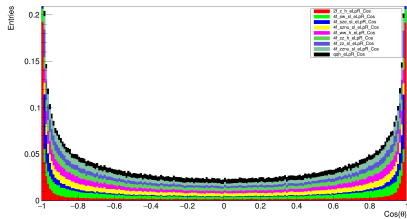
$$\Delta\sigma = \left(\frac{\Delta N_{Total}}{\epsilon L} \right)_{stat} \oplus \left(\frac{\Delta N_B}{\epsilon L} \oplus \sigma \frac{\Delta\epsilon}{\epsilon} \oplus \sigma \frac{\Delta L}{L} \right)_{syst}; \quad \frac{\Delta g_{HZZ}}{g_{HZZ}} = \frac{\Delta\sigma_{ZH}}{2\sigma_{ZH}}$$

- $\Delta N_{Total} = \Delta N_B = 0$ if no fit to the data is performed **Question:**
¿How can the fit be made in the hadronic channel? Next slide.
- $\Delta\epsilon$ is the greatest challenge since it involves the error propagation from the background rejection process.
- $\Delta L/L \sim 0.1\%$ from studies of the ILC and it is negligible compared to the others.

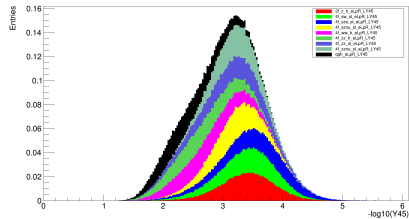
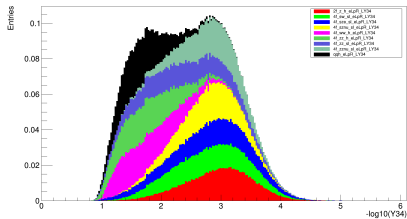
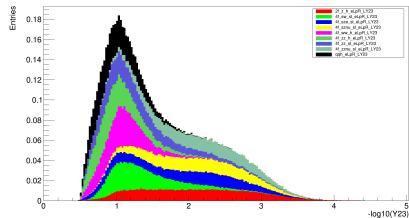
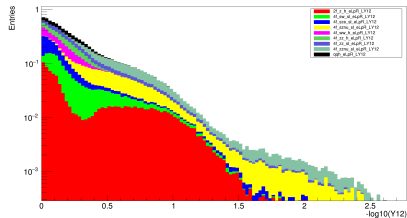
$$F_M(x; M_H, N_S) = N_S \cdot F_S(x; M_H) + N_B \cdot F_B(x)$$



Proposal of background rejection variables: $\text{Cos}(\theta)$, P_t , η and M_Z

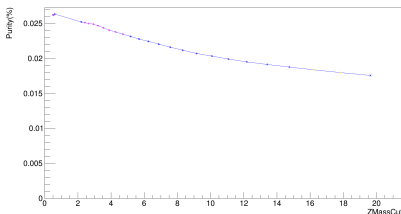
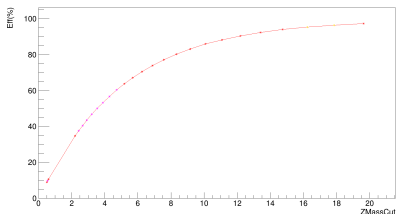


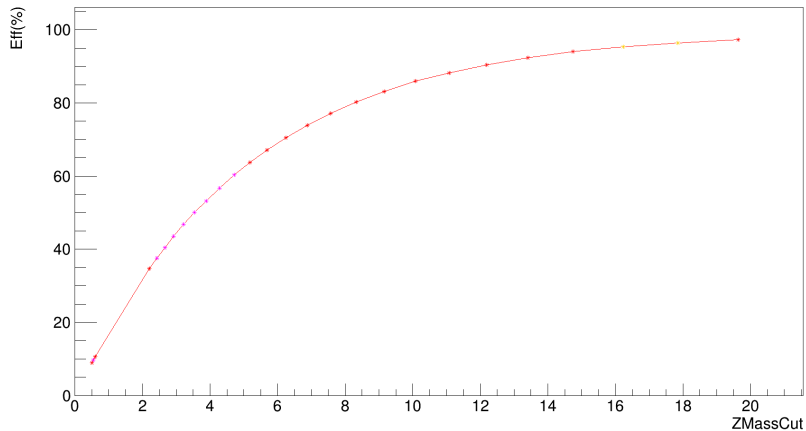
Other background rejection variables: Y_{12} , Y_{23} , Y_{34} and Y_{45}

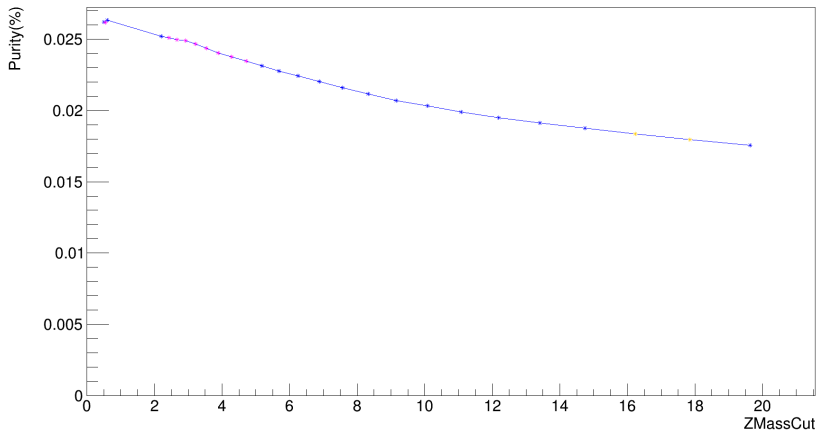


$$\epsilon = \frac{N_{signal}^{sel}}{N_{signal}^{tot}} ; Purity = \frac{N_{signal}^{sel}}{N_{events}^{tot}} \implies \frac{\Delta\sigma}{\sigma} = \frac{\Delta Purity}{Purity} \oplus \frac{\Delta\epsilon}{\epsilon} \oplus \frac{\Delta L}{L}$$

And using as target a resolution of 3% from previous analysis this creates partial and optimal ranges of selection for the cuts. Cut in the Z Mass as an example:







BACKUP

- Processes uncertain if they apply as background: aa_{2f} ; aa_{4f} ;
 ae_{3f} ; ae_{5f}

- The latest MC production *sv02-02-01*, set 250 GeV, reconstructed with the SDHCAL *rv02-02-01.mILD_I5_o2_v02* has now included the overlay events.
- Number of available events 50k in the DST-Merged files analyzed.
- Polarization eLpR (analysis of other polarizations is planned in the future).
- Beam crossing angle correction applied.

The objective of this preliminary analysis is to understand the effect of the newly included overlay into the signal events.

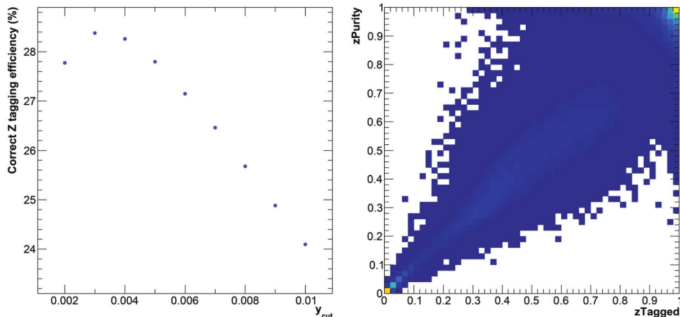
- We run over all possible jet combinatorial possibilities.
- Then we choose the pair that give us the closest mass to the Z known mass.
- Two quality variables are defined:

$$Z_{Tag} = \frac{E_Z^{j_1, j_2}}{E_Z^{Total}} \quad Z_{Purity} = \frac{E_Z^{j_1, j_2}}{E^{j_1, j_2}}$$

where $E_Z^{j_1, j_2}$: MC-truth di-jet energy from the Z, E_Z^{Total} : MC-truth Z energy and E^{j_1, j_2} : MC-truth total di-jet energy.

- The $E_Z^{j_1, j_2}$ is computed through a loop over all the PFOs. The associated MC particle energy is weighted using the RecoMCTrughLink excluding the particle with a Higgs as a parent (pdgId = 25).

The algorithm used in this study was *ee_kt_algorithm* within fast-jet with ExclusiveYCut strategy. An optimization of y_{cut} has to be done in order to maximize the di-jet reconstruction selection efficiency. A di-jet is considered efficient if $Z_{Tag} > 0.9$ and $Z_{Purity} > 0.9$.



The optimal value in this previous study was $y_{cut} = 0.003$.

We want to find the set of parameters from the jet clustering algorithms from which the impact of the overlay is minimal. The two algorithms studied are:

- *Generalized k_t algorithm for e^+e^- collisions (ee_genkt) in the YCut exclusive mode.*

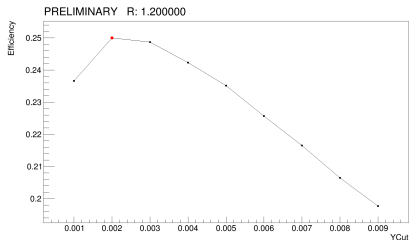
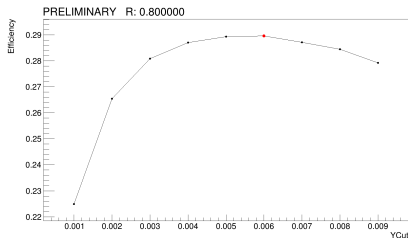
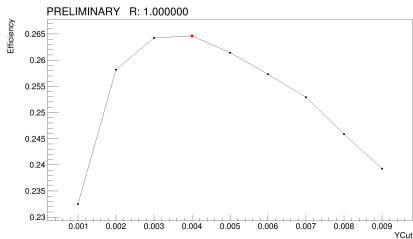
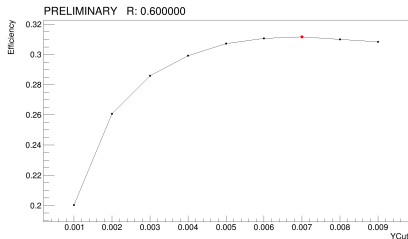
$$d_{ij} = \min(E_i^{2p}, E_j^{2p}) \frac{1 - \cos(\theta_{ij})}{(1 - \cos(R))} \quad ; \quad d_{iB} = E_i^{2p}$$

- *k_t algorithm for e^+e^- collisions (ee_kt) in the YCut exclusive mode.*

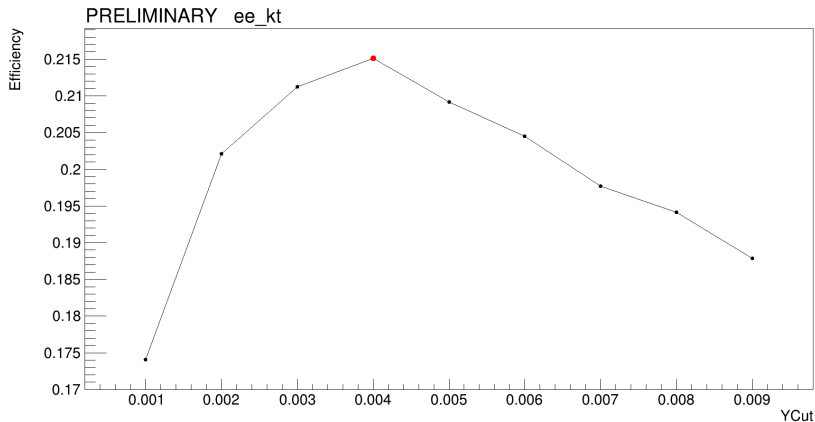
$$d_{ij} = 2\min(E_i^2, E_j^2)(1 - \cos(\theta_{ij}))$$

Parameters: $ee_genkt \rightarrow$ **YCut**, **R** and **P** ; $ee_kt \rightarrow$ **YCut**

YCut scans like the ones in Slide 4 are performed for different values of **R** and **P** for the ee_genkt algorithm and another one for the ee_kt algorithm.



YCut scans like the ones in Slide 4 are performed for different values of **R** and **P** for the ee_genkt algorithm and another one for the ee_kt algorithm.



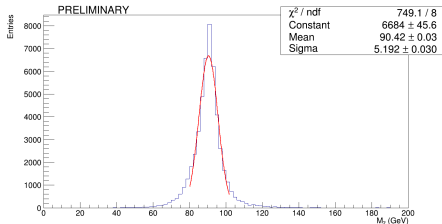
From the studied parameters $R = 0.6$ gives the highest efficiency. However this does not mean that gives the best jet reconstruction. The quality variables computed are:

- Z_{Mass} differences (%). The Z mass is reconstructed, following the previous procedure, with and without overlay. This distribution is fitted to a gauss function and then two differences are computed:

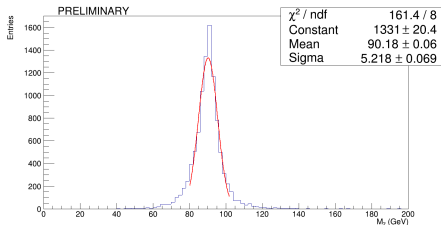
$$ZDiff = |M_Z - M_Z^{NoOverlay}| / M_Z$$

$$OverlayDiff = |M_Z^{Overlay} - M_Z^{NoOverlay}| / M_Z^{NoOverlay}$$

R = 0.6



R = 0.6 No Overlay



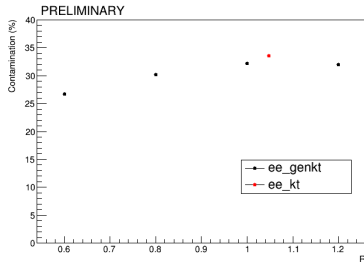
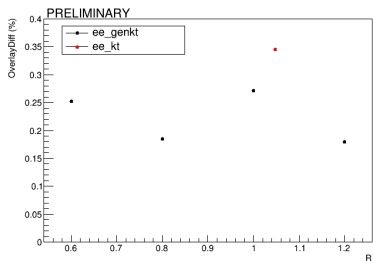
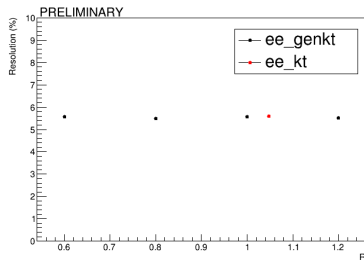
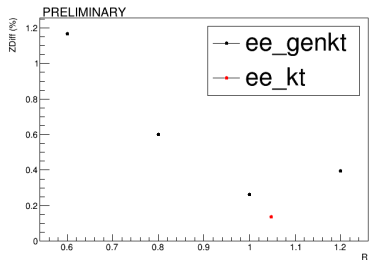
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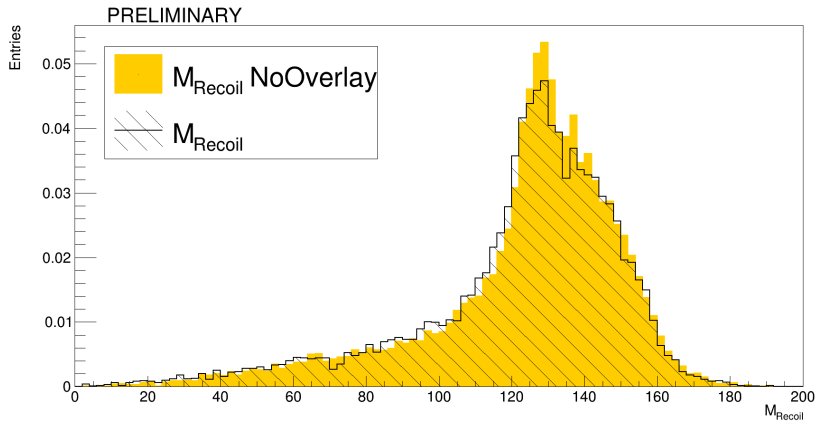
$$ZDiff = |M_Z - M_Z^{NoOverlay}| / M_Z$$

$$OverlayDiff = |M_Z^{Overlay} - M_Z^{NoOverlay}| / M_Z^{NoOverlay}$$

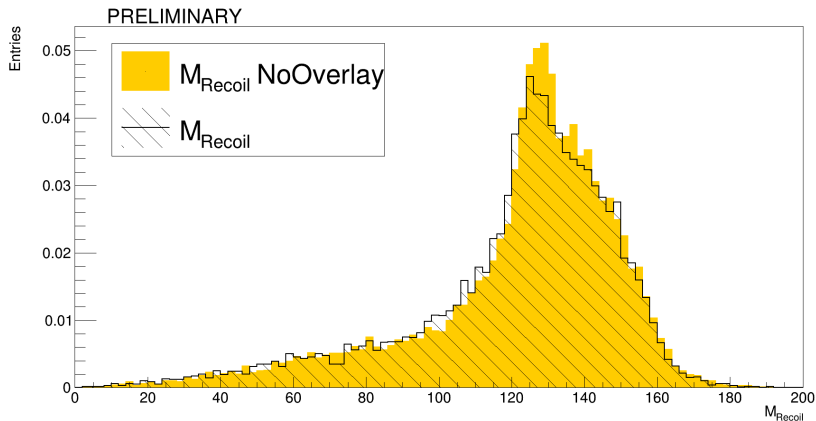
- Z_{Mass} "resolution". (σ / M_Z) from the previous fit.
- Jet Contamination (%). The mean percentage of jets that have some overlay contribution.



With this algorithm the optimal parameters have been found to be $R = 1.0$ and $Y_{\text{Cut}} = 0.004$ and an impact of less than 0.3 % from overlay in the Z reconstruction.



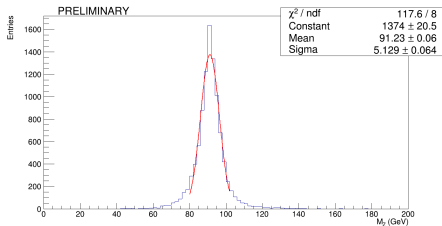
With this algorithm the optimal parameter is $Y_{\text{Cut}} = 0.004$ and an impact of 0.35 % from overlay in the Z reconstruction.



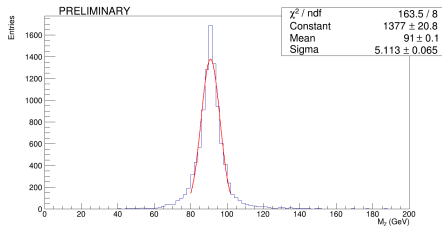
It has been shown the set of parameters that minimizes the overlay effect while efficiently reconstructing the Z mass and showing healthy recoil mass distributions.

A slightly better shape in the recoil mass is obtained from the `ee_kt` method that indicates that it is the algorithm that should be used.

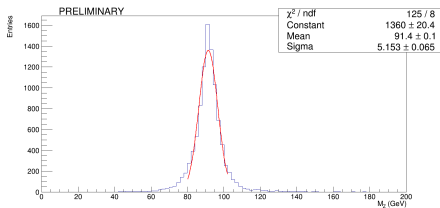
Overlay



No Overlay



Overlay



No Overlay

